

Classical & Quantum Waves

Lecture 13-1

Transport of energy by a wave

• Last lecture, we saw that a traveling wave contains both kinetic and potential energy.

• Total energy is a portion of a vibrating string between $a \leq x \leq b$:

$$E = \frac{1}{2} \mu \int_a^b dx \left[\underbrace{\left(\frac{\partial y}{\partial t} \right)^2}_{\text{Kinetic}} + v^2 \underbrace{\left(\frac{\partial y}{\partial x} \right)^2}_{\text{Potential}} \right]$$

Kinetic
Energy from string
moving up/down in
y-direction

Potential
Energy from
string extending
as wave moves
by

$$\neq \frac{\partial y}{\partial t} \neq v!$$

velocity of wave
is along
x-direction!

• For a pulsed wave:



y-displacement = 0 everywhere except for within the pulse

→ all the energy is inside the pulse (from eq. above)

• For a sinusoidal wave, $y = A \sin(kx - \omega t)$

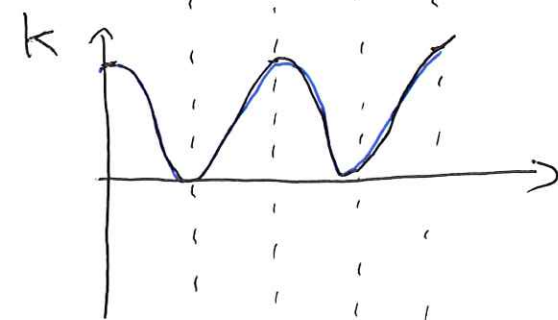
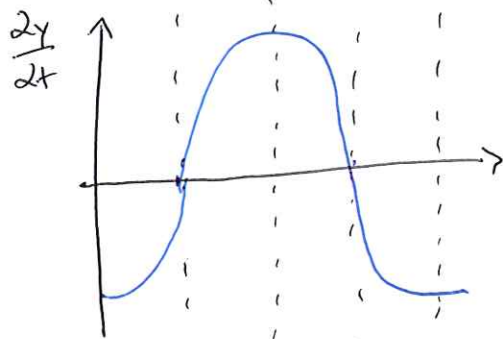
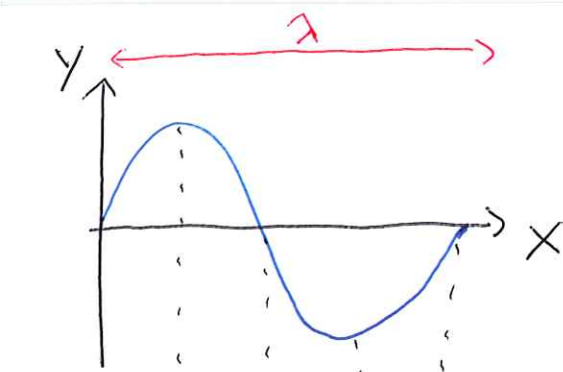
$$\frac{\partial y}{\partial t} = -\omega A \cos(kx - \omega t)$$

• Using above expression for $E = K + U$

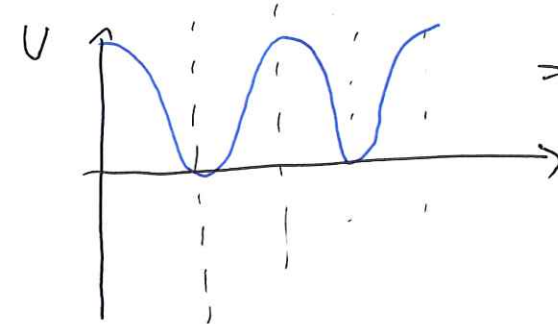
$$K = \frac{1}{2} \mu \delta x \omega^2 A^2 \cos^2(kx - \omega t)$$

$$U = \frac{1}{2} v^2 \mu \delta x k^2 A^2 \cos^2(kx - \omega t)$$

[refer to previous lecture]



$$\Rightarrow \text{max when } \left| \frac{dy}{dt} \right| \text{ max } K \sim \left(\frac{dy}{dt} \right)^2$$



$$\Rightarrow \text{max when } \left| \frac{dy}{dx} \right| \text{ max } U \sim \left(\frac{dy}{dx} \right)^2$$

\hookrightarrow when $y(x)$ passes through zero
(max slope)

\Rightarrow Qualitatively different result from harmonic oscillator!
(analogy w/ coupled HOs & waves breaks down here)

For HO, U was max when displacement max b/c
had $\frac{1}{2}ky^2$ restoring force \rightarrow we don't have
that here.

• Total energy in one λ ^{segment} is $E_{\text{tot}} = \frac{1}{2} \mu \omega^2 A^2 \lambda$ 13-3
[from last lecture]

• The time to travel length λ is $\frac{\lambda}{v}$

• Energy transported per unit time (ie, Power) is then

$$\frac{E_{\text{tot}}}{\frac{\lambda}{v}} = \left[\text{Power} = \frac{1}{2} \mu \omega^2 A^2 v \right] \propto \omega^2, A^2, v$$

//

Waves at discontinuities

• When a wave encounters a discontinuity at a boundary between two different media, in general part of the incoming wave will be reflected, and part will be transmitted.

→ Here we determine the amplitudes and phases of these three waves

• Many examples where this occurs:

• mechanical vibrations

• light / optics

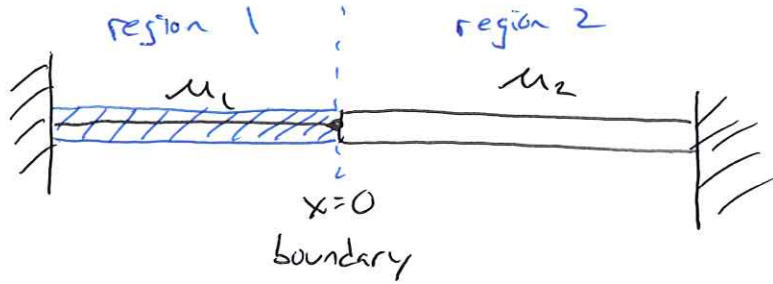
• acoustics

• quantum

...

Consider two strings w/ different mass per unit length 13-41

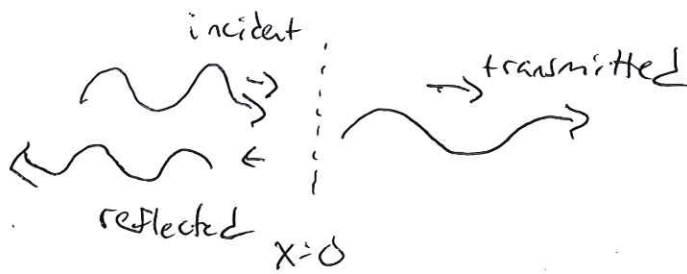
μ_1 and μ_2 . Connected at $x=0$, held under tension T



$$V_1 = \sqrt{\frac{T}{\mu_1}}$$

$$V_2 = \sqrt{\frac{T}{\mu_2}}$$

We generate an ^{incident} wave in region 1, we want to know the amplitude and phase of reflected and transmitted wave at the boundary b/t two regions at $x=0$



$$y_I = A_1 \cos(\omega t - k_1 x)$$

$$y_T = A_2 \cos(\omega t - k_2 x)$$

$$y_R = B_1 \cos(\omega t + k_1 x)$$

\hookrightarrow + $\frac{1}{2}c$ moving to left

Considerations / boundary conditions

1) At the boundary, where two strings connected, displacements at $x=0$ for both strings must be the same

\rightarrow this means that $y_I(x=0) + y_R(x=0) = y_T(x=0)$

region 1 amplitude at $x=0$

region 2 amplitude at $x=0$

Define $y_1 = y_I + y_R$ (waves in region 1)

(waves in region 2) $y_2 = y_T$ (just to keep notation consistent)

\Rightarrow

$$y_1(x=0, t) = y_2(x=0, t)$$

Boundary condition

- This condition also means that the frequencies ω must also be the same on both sides of the boundary

$$\omega_1 = \omega_2 = \omega$$

- However, since wave velocities will be different in two regions, the wavelengths must be different b/c $\lambda = \frac{2\pi}{\omega} v$

$$\Rightarrow \left| k_1 = \frac{\omega}{v_1} \neq k_2 = \omega/v_2 \right|$$

$$\left[\lambda = \frac{2\pi}{k}, \text{ ~~not~~ } \right]$$

$$v = \frac{\omega}{k} = \frac{2\pi f}{\frac{2\pi}{\lambda}} = f\lambda$$

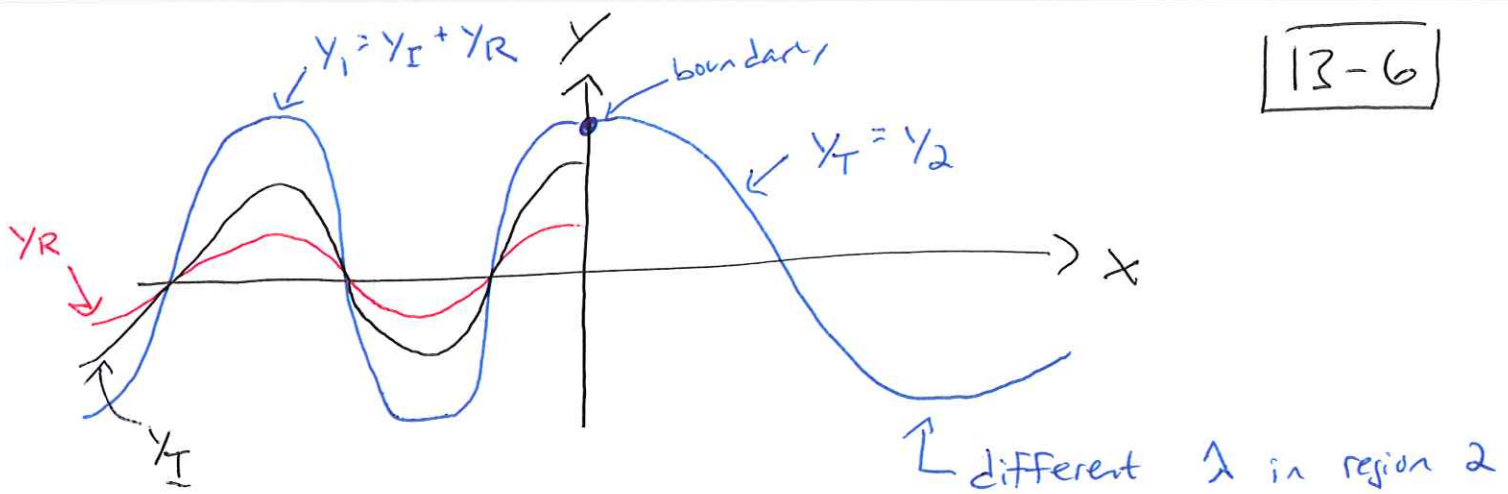
Consideration 2)

- Restoring force at boundary must be continuous, otherwise the finite difference in force acting on an infinitely small segment/mass would lead to infinite acceleration, which is unphysical.

$$\cdot \text{Transverse Force} = T \left(\frac{\partial y}{\partial x} \right) \quad [\text{recall last lecture}]$$

\Rightarrow slopes $\frac{\partial y}{\partial x}$ must be continuous

$$\left| \frac{\partial y_1}{\partial x} (x=0, +) = \frac{\partial y_2}{\partial x} (x=0, +) \right|$$



$$y_1 = y_I + y_R = A_1 \cos(\omega t - k_1 x) + B_1 \cos(\omega t + k_1 x)$$

$$y_2 = y_T = A_2 \cos(\omega t - k_2 x)$$

• Boundary condition (1) gives $y_1 = y_2$ @ $x=0$

$$\Rightarrow A_1 \cos(\omega t - k_1 x) + B_1 \cos(\omega t + k_1 x) = A_2 \cos(\omega t - k_2 x)$$

@ $x=0$. Must be true for all times. Evaluate at $t=0$

$$\Rightarrow \boxed{A_1 + B_1 = A_2}$$

(as in drawing above)

• Boundary condition (2) gives $\frac{\partial y_1}{\partial x} = \frac{\partial y_2}{\partial x}$ @ $x=0$

$$\frac{\partial y_1}{\partial x} = - \underbrace{(-k_1)}_{+k_1} A_1 \sin(\omega t - k_1 x) - k_1 B_1 \sin(\omega t + k_1 x)$$

$$\frac{\partial y_2}{\partial x} = - \underbrace{(-k_2)}_{+k_2} A_2 \sin(\omega t - k_2 x)$$

$$\Rightarrow k_1 A_1 \sin(\omega t - k_1 x) - k_1 B_1 \sin(\omega t + k_1 x) = k_2 A_2 \sin(\omega t - k_2 x) \quad \boxed{13-7}$$

@ $x=0$. Evaluate at $t = \frac{\pi}{2\omega}$ to make the sin term = 1 (we can't do this b/c must be true for all times)

$$\Rightarrow \boxed{k_1 A_1 - k_1 B_1 = k_2 A_2}$$

We want the ratios of amplitudes A_2/A_1 for the transmitted and incident waves and B_1/A_1 for the reflected & incident waves.

$$\rightarrow B_1 = A_2 - A_1 \quad (1)$$

$$(2) \Rightarrow k_1 A_1 - k_1 (A_2 - A_1) = k_2 A_2$$

$$\rightarrow k_1 A_1 + k_1 A_1 - k_1 A_2 = k_2 A_2$$

$$\rightarrow 2k_1 A_1 = A_2 (k_1 + k_2)$$

$$\rightarrow \boxed{\frac{A_2}{A_1} = \frac{2k_1}{k_1 + k_2} = T_{12}}$$

"Transmission coefficient"
(of amplitude)

$$0 < T_{12} < 2$$

[Take limits $k_1 \rightarrow 0$
 $k_2 \rightarrow 0$]

Manipulate same equations to find B_1/A_1

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$$A_2 = B_1 + A_1$$

$$\rightarrow k_1 A_1 - k_1 B_1 = k_2 (B_1 + A_1) \quad (2)$$

$$\rightarrow \cancel{k_1} A_1 (k_1 - k_2) = B_1 (k_1 + k_2)$$

$$\Rightarrow \boxed{\frac{B_1}{A_1} = \frac{k_1 - k_2}{k_1 + k_2} = R_{12}}$$

"Reflection coefficient"
(of amplitude)

take limits
 $-1 < R_{12} < +1$ $\begin{bmatrix} k_1 \rightarrow 0 \\ k_2 \rightarrow 0 \end{bmatrix}$

One can also show, by manipulating the ~~two~~ ^{transmission}

8 reflection coeff. equations, that

$$\boxed{T_{12} = 1 + R_{12}}$$

[to show this, solve for $(k_1 + k_2)$,
plug in the other eq. etc.]

~~From this, ^{one} ~~we~~ can show that T_{12} is always positive
and between 0 and 2~~

~~R_{12} is between -1 and +1~~

$$R_{12} = \frac{k_1 - k_2}{k_1 + k_2}$$

shows that the sign of

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the reflected wave depends on whether k_2 is less than or greater than k_1 .

If $k_2 < k_1$, $R_{12} = \frac{B_1}{A_1}$ is positive & reflected wave is in phase w/ incident wave (see earlier drawing)

$$\lambda_2 > \lambda_1$$

$$v_2 > v_1$$

$$\mu_2 < \mu_1$$

If $k_2 > k_1$, $R_{12} = \frac{B_1}{A_1}$ is negative & reflected wave is ~~the~~ opposite sign (π out of phase) compared to incident wave

$T_{12} = \frac{A_2}{A_1}$ shows that transmitted wave always positive and in phase with incident wave

Rewrite T_{12} & R_{12} in terms of μ

$$v = \frac{\omega}{k} = \sqrt{\frac{T}{\mu}} \Rightarrow k = \omega \sqrt{\frac{\mu}{T}}$$

$$\Rightarrow T_{12} = \frac{2k_1}{k_1 + k_2} = \frac{2\sqrt{\mu_1}}{\sqrt{\mu_1} + \sqrt{\mu_2}} \quad [\text{other terms cancel}]$$

$$R_{12} = \frac{k_1 - k_2}{k_1 + k_2} = \frac{\sqrt{\mu_1} - \sqrt{\mu_2}}{\sqrt{\mu_1} + \sqrt{\mu_2}}$$

~~As~~ As $\mu_2 \rightarrow \infty$, region 2 becomes a rigid wall.

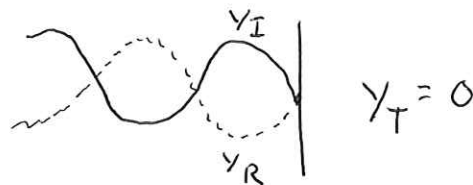
$$\cancel{R_{12}} \quad T_{12} \rightarrow 0$$

no transmitted
wave

$$\text{recall } T_{12} = 1 + R_{12}$$

$$\Rightarrow R_{12} = -1$$

→ reflected wave at rigid wall reverses amplitude by getting π phase shift



→ This concept is universal!

- light / EM waves
- sound waves
- Quantum waves

Waves in two- & three-dimensions

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- Description of waves in one dimension can be extended to 2D and 3D

2D wave eq:
$$\frac{\partial^2 z}{\partial x^2} + \frac{\partial^2 z}{\partial y^2} = \frac{1}{v^2} \frac{\partial^2 z}{\partial t^2}$$

displacement in
z-direction

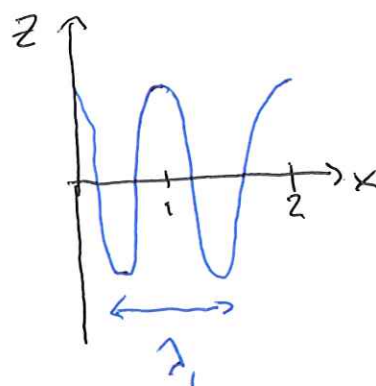
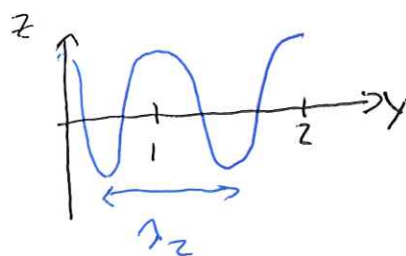
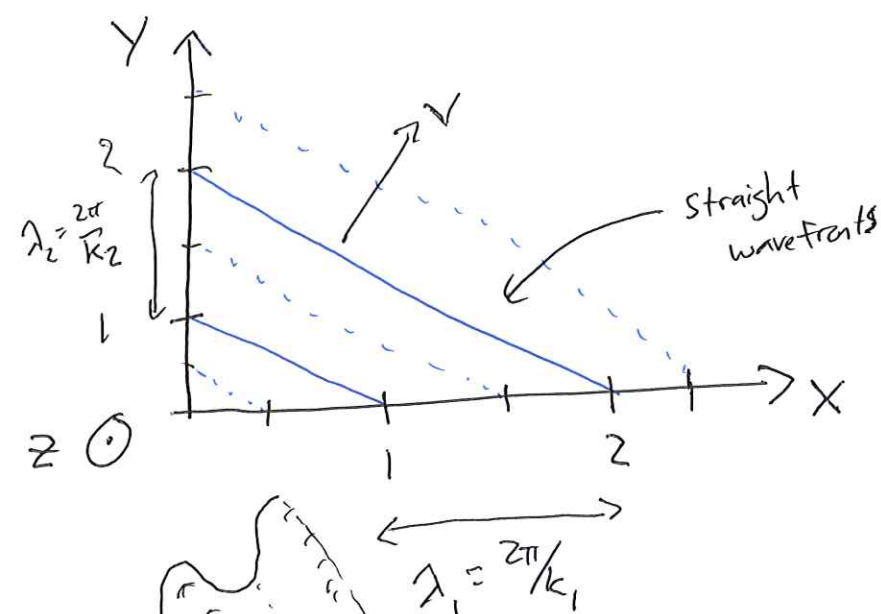
Solution:
$$z(x, y, t) = A \cos(k_1 x + k_2 y - \omega t)$$

propagating in
x, y direction

velocity:
$$v = \frac{\omega}{\sqrt{k_1^2 + k_2^2}} = \frac{\omega}{k}$$

$$k = \sqrt{k_1^2 + k_2^2}$$

Snapshot of wave fronts:



[Attempted 3D illustration...]

2D w/ circular symmetry

13-12

$$\boxed{\frac{\partial^2 z}{\partial r^2} + \frac{1}{r} \frac{\partial z}{\partial r} = \frac{1}{v^2} \frac{\partial^2 z}{\partial t^2}}$$

(no proof)

[pebble in a pond]



Solutions to this are called Bessel functions

→ for $r \rightarrow \infty$:

$$\boxed{\frac{\partial^2 z}{\partial r^2} = \frac{1}{v^2} \frac{\partial^2 z}{\partial t^2}}$$

→ solved by 1D sine wave

$$z(r, t) = A \cos(kr - \omega t)$$

3-dimensions

$$\boxed{\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} = \frac{1}{v^2} \frac{\partial^2 \psi}{\partial t^2}}$$

$= \Delta \psi$ w/ $\Delta =$ Laplace operator

$$A \cos(k_1 x + k_2 y + k_3 z - \omega t)$$

Solution: 3D sine wave ~~$\frac{\partial^2 \psi}{\partial x^2}$~~ $\psi(x, y, z, t) = \cancel{A \cos(k_1 x + k_2 y + k_3 z - \omega t)}$

$$v = \frac{\omega}{k} \quad \text{w/} \quad k = \sqrt{k_1^2 + k_2^2 + k_3^2}$$

3D (spherical symmetry)

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3D spherical wave eq:
$$\frac{\partial^2 \psi}{\partial r^2} + \frac{2}{r} \frac{\partial \psi}{\partial r} = \frac{1}{v^2} \frac{\partial^2 \psi}{\partial t^2}$$

Solution: 3D spherical
waves

$$\psi(r, t) = \frac{A}{r} \cos(\omega t - kr)$$

Amplitude decreases as $\frac{1}{r}$

$$\rightarrow \text{Energy} \sim \text{amplitude}^2 \propto \frac{1}{r^2}$$

• Higher dimension waves relevant to

acoustic, EM, quantum... waves!